

Skyrme Crystals

Martin Speight

Joint work with Derek Harland and Paul Leask

14/3/23

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- Skyrme energy

$$E(\varphi) = \int_M |d\varphi|^2 + \frac{1}{4} |\varphi^* \omega|^2 + V(\varphi)$$

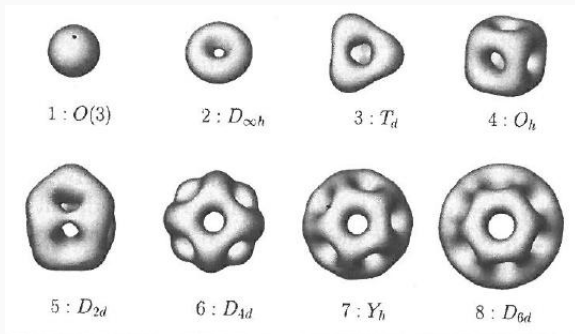
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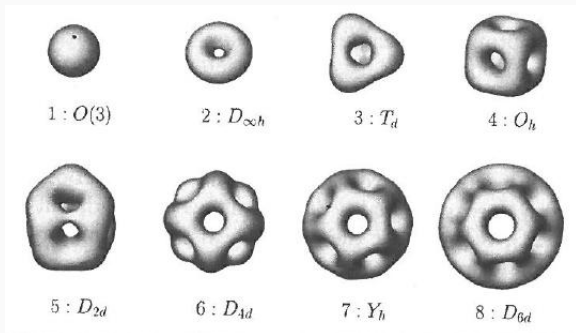
- Faddeev bound: $E(\varphi) \geq E_0 |B|$, unattainable
- Degree B minimizer \leftrightarrow nucleus of atomic weight B

- Numerics



Battye and Sutcliffe

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Battye and Sutcliffe

- E/BE_0 monotonically decreases e.g. 1.232 ($B = 1$), 1.096 ($B = 8$).

- Suggests Skyrmions may be able to form a **crystal**

$$\varphi : \mathbb{R}^3/\Lambda \rightarrow G, \quad \Lambda = \{n_1\mathbf{X}_1 + n_2\mathbf{X}_2 + n_3\mathbf{X}_3 : \mathbf{n} \in \mathbb{Z}^3\}$$

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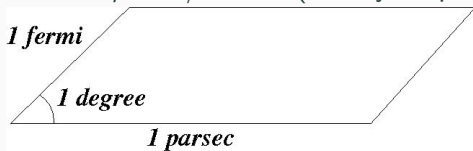
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- But is this really a crystal? Given **any** Λ , B , there exists a degree B minimizer $\varphi : \mathbb{R}^3/\Lambda \rightarrow G$ (Auckly, Kapitanski).



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For most Λ , lifted map $\mathbb{R}^3 \rightarrow G$ clearly isn't a genuine solution: artifact of bc's.

- Given a minimizer $\varphi : \mathbb{R}^k/\Lambda \rightarrow N$ of some energy functional $E(\varphi)$, when is the lifted map $\mathbb{R}^k \rightarrow N$ a genuine crystal?

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- Should be critical (in fact stable) with respect to variations of Λ too.

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- Identify them all with $M = \mathbb{R}^3/\mathbb{Z}^3$, the cubic torus. Now mfd is fixed, but **metric** depends on Λ

$$g_\Lambda = g_{ij} dx_i dx_j, \quad g_{ij} = \mathbf{X}_i \cdot \mathbf{X}_j \text{ const}$$

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- So now

$$E(\varphi, g) = \int_{T^3} (|d\varphi|_g^2 + \frac{1}{4} |\varphi^* \omega|_g^2 + V(\varphi)) \text{vol}_g$$

and we want to minimize w.r.t. both $\varphi \in C_B^2(T^3, G)$ **and** $g \in SPD_3$ (space of symmetric positive definite 3×3 matrices)

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- Fix g : $E(\varphi)$ certainly attains a min in each homotopy class (at least in H^1 - low regularity) – Auckly, Kapitanski
- What if we fix φ ? Does $E(g)$ attain a min?

Existence and uniqueness of minimizing metrics

- Want to think of E , for a **fixed** $\varphi : T^3 = \mathbb{R}^3/\mathbb{Z}^3 \rightarrow G$ as a function of the metric g on T^3 :

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- **Theorem** Let $\varphi : T^3 \rightarrow G$ be C^1 and somewhere immersive. Then E_φ attains a global minimum at some $g_* \in SPD_3$ and has no other critical points.
- Proof: First note that

$$E_\varphi(g) = \sqrt{|g|} \operatorname{tr}(Hg^{-1}) + \frac{1}{\sqrt{|g|}} \operatorname{tr}(\Omega g) + C\sqrt{|g|}$$

where $H, \Omega \in SPD_3$ and $C \in [0, \infty)$ are fixed.

Existence and uniqueness of minimizing metrics

$$E_2(g) = \int_{T^3} \varphi^* h(\partial_i, \partial_j) g^{ij} \sqrt{|g|} d^3x = \sqrt{|g|} g_{ij}^{-1} H_{ij}$$

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$$E_4(g) = \frac{1}{4} \|\varphi^* \omega\|_{L^2(g)}^2 = \frac{1}{4} \int_{T^3} \frac{1}{|g|} g(X_\varphi, X_\varphi) \text{vol}_g = \frac{g_{ij}}{\sqrt{|g|}} \Omega_{ij}$$

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$$(\lambda, \mathcal{O}) \mapsto \mathcal{O} \begin{pmatrix} \lambda_1 & 0 & 0 \\ 0 & \lambda_2 & 0 \\ 0 & 0 & \lambda_3 \end{pmatrix} \mathcal{O}^T = \mathcal{O} D_\lambda \mathcal{O}^T$$

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J.M. Speight (University of Leeds) • We will show $E \circ f : (0, \infty)^3 \times O(3) \rightarrow \mathbb{R}$ attains a min

$$\begin{aligned}(E \circ f)(\lambda, \theta) &= \operatorname{tr}(H(\theta D_\lambda \theta^{-1})^{-1}) + \operatorname{tr}(\Omega \theta D_\lambda \theta^{-1}) + \frac{C}{\lambda_1 \lambda_2 \lambda_3} \\ &= \operatorname{tr}(\theta^{-1} H \theta D_\lambda^{-1}) + \operatorname{tr}(\theta^{-1} \Omega \theta D_\lambda) + \frac{C}{\lambda_1 \lambda_2 \lambda_3}\end{aligned}$$

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- Consider the smooth functions $O(3) \rightarrow (0, \infty)$

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- Exists $\alpha > 0$ s.t. for all $(\boldsymbol{\lambda}, \mathcal{O})$,

$$(E \circ f)(\boldsymbol{\lambda}, \mathcal{O}) \geq \alpha \left(\frac{1}{\lambda_1} + \frac{1}{\lambda_2} + \frac{1}{\lambda_3} + \lambda_1 + \lambda_2 + \lambda_3 \right). \quad (*)$$

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- Continuity of E implies $(E \circ f)(\lambda_*, \mathcal{O}_*) = E_*$, i.e. $E \circ f$ attains a min
- Let $\mathcal{G}_* = \mathcal{O}_* D_{\lambda_*} \mathcal{O}_*^{-1} \in SPD_3$. E_φ attains a min at $\mathbf{g}_* = \mathcal{G}_* / \det(\mathcal{G}_*)$.

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- Metric on SPD_3 ?

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- Geodesic through $\mathcal{G}(0)$: $\mathcal{G}(t) = A \exp(t\xi) A^T$ where $AA^T = \mathcal{G}(0)$

$$E_4(\mathcal{G}) = \text{tr}(\Omega \mathcal{G})$$

$$\begin{aligned} E_4(\mathcal{G}(t)) &= \text{tr}(\Omega A \exp(t\xi) A^T) \\ &= \text{tr}(\Omega_A \exp(t\xi)), \quad \Omega_A = A^T \Omega A \end{aligned}$$

$$\left. \frac{d^2}{dt^2} E_4(\mathcal{G}(t)) \right|_{t=0} = \text{tr}(\Omega_A \xi^2) > 0$$

- So E_4 is strictly convex.

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- Hence $E_0 = \det \circ \iota$ is strictly convex
- So $E = E_2 + E_4 + E_0$ is strictly convex. Hence it has at most one critical point. (Assume \mathcal{G}_* , \mathcal{G}_{**} both cps, apply Rolle's Theorem to $(E \circ \gamma)'$ where γ is the geodesic between them.)

- Minimize $E : C^2(T^3, S^3) \times SPD_3 \rightarrow \mathbb{R}$

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- Newton flow: pick $x(0) \in M$, solve

$$\ddot{x} = -(\text{grad } E)(x)$$

with $\dot{x}(0) = 0$.

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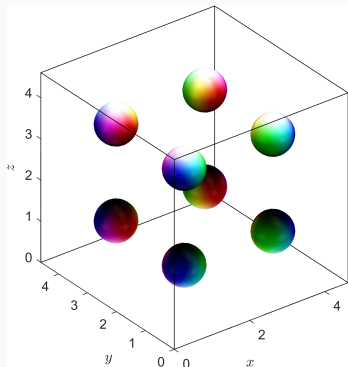
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- Converges much faster than gradient flow.

The Kugler-Shtrikman crystal (massless model)

$$E = \|d\varphi\|^2 + \frac{1}{4}\|\varphi^*\omega\|^2$$



$$\begin{aligned}(x_1, x_2, x_3) &\mapsto (x_2, x_3, x_1) \\ (\varphi_0, \varphi_1, \varphi_2, \varphi_3) &\mapsto (\varphi_0, \varphi_2, \varphi_3, \varphi_1)\end{aligned}$$

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- No reason to expect **degenerate** critical points to survive perturbation

An instructive toy model

$$E_t : \mathbb{R}^2 \rightarrow \mathbb{R}, \quad E_t(x, y) = x^2 + ty^2$$

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- Also a solution of $\nabla E_t = 0$ by PSC

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- E_0 invariant under action of $G_0 = SO(4) \times Aut(T^3)$

The case of the KS crystal

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- Then $M^{\Gamma_R} = M$ and $(R, e) \cdot \varphi_{KS}$ is certainly **not** a nondegenerate cp of E_0

The case of the KS crystal

- For

$$R \in \left\{ \mathbb{I}_4, \underbrace{\left(\begin{array}{c} (0, 1, 1, 1)/\sqrt{3} \\ * \end{array} \right)}_{R_\alpha}, \underbrace{\left(\begin{array}{c} (0, 0, 0, 1) \\ * \end{array} \right)}_{R_{sheet}}, \underbrace{\left(\begin{array}{c} (0, 0, 1, 1)/\sqrt{2} \\ * \end{array} \right)}_{R_{chain}} \right\}$$

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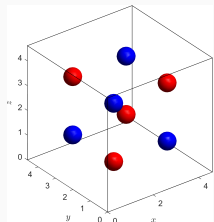
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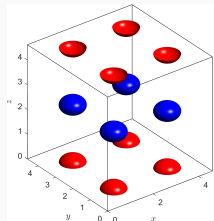
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- PSC implies $\varphi(t)$ also a cp of E_t .

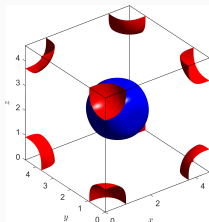
The KS crystals that (should) survive



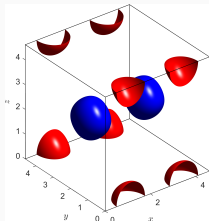
φ_{KS}



$R_{sheet}\varphi_{KS}$



$R_{\alpha}\varphi_{KS}$

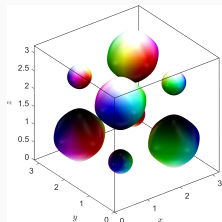


$R_{chain}\varphi_{KS}$

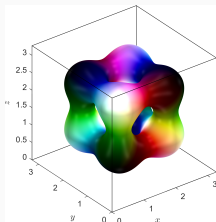
$$\varphi_0 = 0.9$$

$$\varphi_0 = -0.9$$

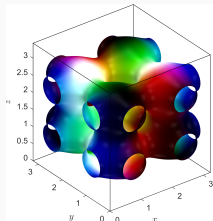
Skyrme crystals at pion mass $t = 1$



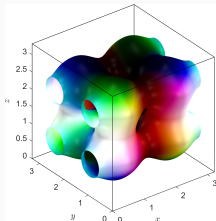
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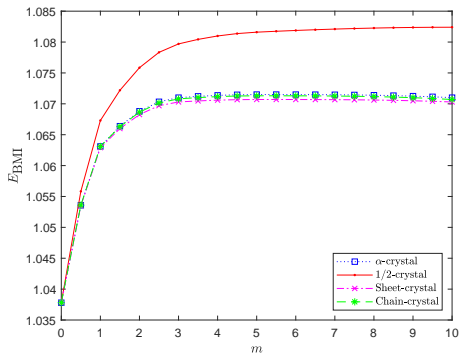


$R_{sheet}\varphi_{KS}$



$R_{chain}\varphi_{KS}$

Energy ordering: sheet < chain < α < KS



$$g_{sheet} = \begin{pmatrix} L_1 & 0 & 0 \\ 0 & L_1 & 0 \\ 0 & 0 & L_3 \end{pmatrix}$$

$$L_3 > L_1$$

$$g_{chain} = \begin{pmatrix} L_1 & 0 & 0 \\ 0 & L_2 & 0 \\ 0 & 0 & L_2 \end{pmatrix}$$

$$L_2 > L_1$$

trigonal, but **not** cubic!

$$U_{KS} = \begin{pmatrix} 165.2 & 0 & 0 \\ 0 & 165.2 & 0 \\ 0 & 0 & 165.2 \end{pmatrix}, \quad U_{\alpha} = \begin{pmatrix} 135.5 & 0 & 0 \\ 0 & 135.5 & 0 \\ 0 & 0 & 167.3 \end{pmatrix},$$
$$U_{sheet} = \begin{pmatrix} 135.8 & 0 & 0 \\ 0 & 135.8 & 0 \\ 0 & 0 & 166.8 \end{pmatrix}, \quad U_{chain} = \begin{pmatrix} 135.6 & 0 & 0 \\ 0 & 135.7 & 0 \\ 0 & 0 & 167.2 \end{pmatrix}.$$

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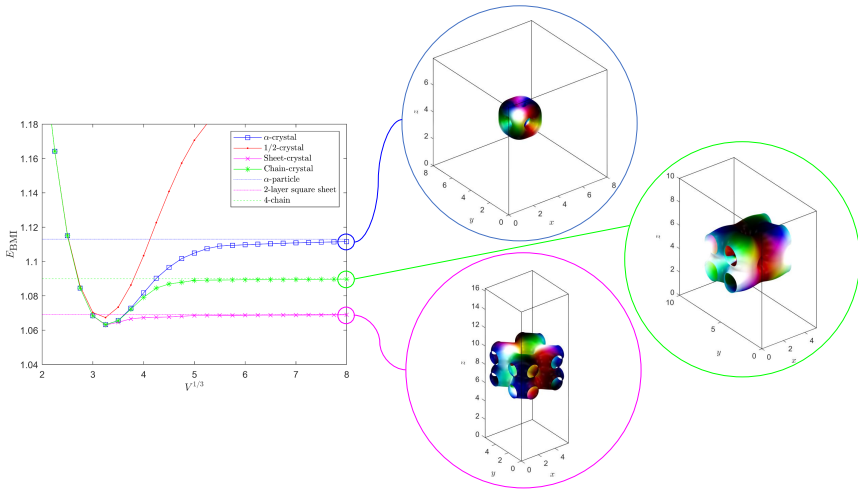
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- Can again solve numerically by ANF

Optimal crystals at fixed baryon density



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Concluding remarks

- Energetically optimal soliton lattices do **not** necessarily have cubic (or triangular) symmetry!
- Many examples in condensed matter. True also for nuclear Skyrme model with massive pions
- Extreme case: baby Skyrme model $\varphi : M^2 \rightarrow S^2$

$$E(\varphi) = \int_M \left(\frac{1}{2} |d\varphi|^2 + \frac{1}{2} |\varphi^* \omega|^2 + V(\varphi) \right).$$

Given **any** period lattice $\Lambda \subset \mathbb{R}^2$, can cook up a smooth potential $V : S^2 \rightarrow [0, \infty)$ s.t. $E(\varphi, g)$ has a global min at (φ_*, g_Λ) with φ_* degree 2 and holomorphic.

- So this crazy lattice **is** the period lattice of a baby Skyrmion crystal, at least for a (highly contrived) choice of V !

