#### The Geometry of Soliton Moduli Spaces

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• Smooth, spatially localized, lump-like solutions of relativistic nonlinear wave equations

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- Like strings...

hypothetical particles, resolve many theoretical puzzles in HEP

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#### ...only better!

exist in real world: magnetic flux tubes in superconductors, magnetic bubbles, optical pulses, crystal dislocations

Interesting special case: static solitons exert no net force on each other



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• Moduli space of static *n*-soliton solutions  $M_n$ , dim  $M_n = n \dim M_1$ 

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Low energy dynamics reduces to geodesic motion in M<sub>n</sub>!

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- Low energy dynamics reduces to geodesic motion in M<sub>n</sub>!
- Soliton dynamics ←→ Riemannian geometry

• Planar antiferromagnets  $\to \mathbb{C}P^1$  model

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- The Bogomol'nyi argument, M<sub>n</sub>
- The metric on M<sub>n</sub>, soliton scattering
- Other solitons
- Open problems

• Square spin lattice:  $\mathbf{S} : \mathbb{Z} \times \mathbb{Z} \to S^2$ 

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First order, spin couples to nearest neighbours.

- Square spin lattice:  $\mathbf{S} : \mathbb{Z} \times \mathbb{Z} \to S^2$
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First order, spin couples to nearest neighbours.

Continuum limit?



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$$\begin{array}{ll} \displaystyle \frac{d\mathbf{A}_{\alpha\beta}}{d\tau} & = & -(\mathbf{B}_{\alpha,\beta-1} + \mathbf{B}_{\alpha\beta} + \mathbf{B}_{\alpha-1,\beta} + \mathbf{B}_{\alpha-1,\beta-1}) \\ \displaystyle \frac{d\mathbf{B}_{\alpha\beta}}{d\tau} & = & -(\mathbf{A}_{\alpha+1,\beta} + \mathbf{A}_{\alpha+1,\beta+1} + \mathbf{A}_{\alpha,\beta+1} + \mathbf{A}_{\alpha,\beta}) \end{array}$$

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- $x = \alpha \delta$ ,  $y = \beta \delta$ ,  $t = 2\tau \delta$
- Assumption:

$$\begin{array}{c} \mathbf{A}_{\alpha,\beta} \\ \mathbf{B}_{\alpha,\beta} \end{array} \right\} \quad \stackrel{\delta \to 0}{\longrightarrow} \quad \left\{ \begin{array}{c} \mathsf{A}(x,y) \\ \mathsf{B}(x,y) \end{array} \right.$$

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• 
$$x = \alpha \delta$$
,  $y = \beta \delta$ ,  $t = 2\tau \delta$ 

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• Replace  $\mathbf{A}_{\alpha+1,\beta}$  by  $\mathbf{A} + \delta \mathbf{A}_x + \frac{1}{2} \delta^2 \mathbf{A}_{xx} + \cdots$  etc

• Work to order  $\delta^2$ 

$$\begin{aligned} 2\delta A_t &= -A \times [4B - 2\delta(B_x + B_y) + \delta^2(B_{xx} + B_{yy} + B_{xy})] \\ 2\delta B_t &= -B \times [4A + 2\delta(A_x + A_y) + \delta^2(A_{xx} + A_{yy} + A_{xy})] \end{aligned}$$

$$2\delta A_t = -A \times [4B - 2\delta(B_x + B_y) + \delta^2(B_{xx} + B_{yy} + B_{xy})]$$
  
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• New fields: 
$$\mathbf{m} = \frac{1}{2}(A+B)$$
  $\phi = \frac{1}{2}(A-B)$ 

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 $\bullet \ |\textbf{m}| = \textit{O}(\delta) \quad , \quad |\phi| = 1 + \textit{O}(\delta^2)$ 

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$$\mathbf{m} = \frac{1}{2}(A+B)$$
  $\phi = \frac{1}{2}(A-B)$   
•  $|\mathbf{m}| = O(\delta)$  ,  $|\phi| = 1 + O(\delta^2)$ 

$$\mathbf{m}_{t} = -(\partial_{x} + \partial_{y})[\mathbf{m} \times \boldsymbol{\varphi}] + \frac{\delta}{4}[2\boldsymbol{\varphi} \times (\boldsymbol{\varphi}_{xx} + \boldsymbol{\varphi}_{yy} + \boldsymbol{\varphi}_{xy}) \quad (1)$$
  
$$\delta \boldsymbol{\varphi}_{t} = 4\mathbf{m} \times \boldsymbol{\varphi} - \delta \boldsymbol{\varphi} \times (\boldsymbol{\varphi}_{x} + \boldsymbol{\varphi}_{y}) + O(\delta^{2}) \quad (2)$$

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• Solve (2): 
$$\mathbf{m} = \frac{\delta}{4} \big[ \mathbf{\varphi} \times \mathbf{\varphi}_t - \mathbf{\varphi}_x - \mathbf{\varphi}_y \big] + O(\delta^2)$$

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• Subst in (1):  $\phi \times \phi_{tt} = \phi \times (\phi_{xx} + \phi_{yy}) + O(\delta)$ 

• Leading order

$$\mathbf{\phi} \times \Box \mathbf{\phi} = \mathbf{\phi} \times (\mathbf{\phi}_{tt} - \mathbf{\phi}_{xx} - \mathbf{\phi}_{yy}) = \mathbf{0}$$

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$$\begin{split} \phi \times \Box \phi &= \phi \times \left( \phi_{tt} - \phi_{xx} - \phi_{yy} \right) = 0 \\ \Box \phi - \left( \phi \cdot \Box \phi \right) \phi &= 0 \quad (*) \end{split}$$

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Nonlinear wave equation! Lorentz invariant!

• Variational formulation: action of field  $\phi:\mathbb{R}\times\Sigma\to S^2$ 

$$S[\varphi] = \frac{1}{2} \int_{\mathbb{R} \times \Sigma} \left( |\varphi_t|^2 - |\varphi_x|^2 - |\varphi_y|^2 \right) dt \, dx \, dy$$

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 $\varphi$  solves (\*) iff  $\varphi$  a critical point of S.

#### Leading order

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• Variational formulation: action of field  $\phi:\mathbb{R}\times\Sigma\to \mathit{S}^2$ 

$$S[\varphi] = \frac{1}{2} \int_{\mathbb{R} \times \Sigma} \left( |\varphi_t|^2 - |\varphi_x|^2 - |\varphi_y|^2 \right) dt \, dx \, dy$$

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 $\varphi$  solves (\*) iff  $\varphi$  a critical point of S.

• Physicists call this the  $\mathbb{C}P^1$  model

 $\varphi_{tt} - \Delta \varphi - [\varphi \cdot (\varphi_{tt} - \Delta \varphi)] \varphi = 0$ 

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Static solutions are critical points of potential energy

$$E = \frac{1}{2} \int_{\Sigma} |\varphi_x|^2 + |\varphi_y|^2$$

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- Fields fall into disjoint homotopy classes labelled by  $n = \deg \phi \in \mathbb{Z}$ . WLOG can assume  $n \ge 0$
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$$E = 4\pi n \quad \Leftrightarrow \quad \phi_x + \phi \times \phi_y = 0 \quad \text{1st order PDE!}$$

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$$J: T_{\phi}S^2 \rightarrow T_{\phi}S^2$$
 s.t.  $J^2 = -1$ 

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- Bogomol'nyi equation equivalent to dφ ∘ J<sub>Σ</sub> = J ∘ dφ that is, φ : Σ → S<sup>2</sup> is holomorphic

### Moduli space





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$$u(z) = \frac{a_0 + a_1 z + \dots + a_n z^n}{b_0 + b_1 z + \dots + b_n z^n}$$

• Boundary condition:  $\varphi(\infty) = (0, 0, 1) \Rightarrow b_n = 0.$ 

$$u(z) = \frac{a_0 + a_1 z + \dots + a_{n-1} z^{n-1} + z^n}{b_0 + b_1 z + \dots + b_{n-1} z^{n-1}}$$

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• Moduli space  $M_n = \operatorname{Rat}_n^* \subset \mathbb{C}^{2n}$ , open



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• 
$$M_1 = \mathbb{C} \times \mathbb{C}^{\times}$$
  
 $u(z) = \frac{a_0 + z}{b_0}$ 

• Position  $-a_0$ , width  $|b_0|$ , orientation  $arg(b_0)$ 



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### Moduli space

- $M_2$ , complicated manifold  $\dim_{\mathbb{C}} M_2 = 4$ .
- Expect energy to localize around zeros of *u*. OK if well-separated
- Lose identity when close, e.g.  $u = z^2$





$$S = \int dt \left\{ \frac{1}{2} \left( \int_{\Sigma} |\varphi_t|^2 \right) - E(\varphi) \right\}$$

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**Geodesic** motion on  $(M_n, \gamma)$  where  $\gamma = L^2$  metric.

# $L^2$ metric



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•  $X \in T_{\phi}M_n$  is a section of  $\phi^{-1}TS^2$ 

# L<sup>2</sup> metric



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 Geodesic motion is constant speed motion along "straightest possible" curve

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In general, geodesic flow very complicated

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Lumps of equal width  $\sim \lambda^{-\frac{1}{2}}$  located where  $z^2 = -\mu$ 





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Hermitian  $\Rightarrow \omega(Y, X) = -\omega(X, Y)$ 

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Hermitian  $\Rightarrow \omega(Y, X) = -\omega(X, Y)$ Kähler  $\Rightarrow d\omega = 0$ 

- $M_n \subset \mathbb{C}^{2n}$  is itself a **complex** manifold, has a  $J : T_{\phi}M_n \to T_{\phi}M_n$ ,  $J^2 = -1$
- Two natural structures on each tangent space, *J* and γ. Are they compatible?
- Yes:  $\gamma(JX, JY) = \gamma(X, Y)$  for all  $X, Y \in T_{\phi}M_n$  $\gamma$  is Hermitian
- Even better,  $\nabla J = 0$  (*J* invariant under parallel transport)  $\gamma$  is Kähler
- Alternative characterization: define K\u00e4hler form

 $\omega(X,Y) = \gamma(JX,Y)$ 

Hermitian  $\Rightarrow \omega(Y, X) = -\omega(X, Y)$ Kähler  $\Rightarrow d\omega = 0$ 

• Consequence: centre of mass motion decouples,  $M_n = \mathbb{C} \times M_n^{\text{red}}$ 

$$\gamma = 4\pi n \gamma_{\rm Euc} + \gamma_{\rm red}$$

• Integer topological charge n

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- Topological energy bound E ≥ E<sub>0</sub>n Attained by solutions of a first order nonlinear PDE system, "holomorphic"

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- 90° head-on scattering

• Obvious generalization:

$$\phi: \Sigma \rightarrow N, \qquad E = rac{1}{2} \int_{\Sigma} |\mathrm{d} \phi|^2$$

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 $\Sigma$ , *N* Riemannian mfds (harmonic map problem)

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 $\Sigma$ , *N* Riemannian mfds (harmonic map problem)

- Antiferromagnets:  $\Sigma = \mathbb{C}$ ,  $N = S^2$
- Inhomogeneous antiferromagnets:  $\Sigma = \Sigma^2$ ,  $N = S^2$

Obvious generalization:

$$\phi: \Sigma \to \textit{N}, \qquad \textit{E} = \frac{1}{2} \int_{\Sigma} |\mathrm{d}\phi|^2$$

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- Antiferromagnets:  $\Sigma = \mathbb{C}$ ,  $N = S^2$
- Inhomogeneous antiferromagnets:  $\Sigma = \Sigma^2$ ,  $N = S^2$
- $\Sigma$ , *N* Kähler, keeps key features:  $n = [\phi^* \omega]$  $M_n = \text{hol}_n(\Sigma, N)$ , Kähler

- $\phi$  a section of a complex vector bundle W over  $\Sigma$
- $\nabla$  = unitary connexion on *W*, curvature *F*

$$E = \frac{1}{2} \int_{\Sigma} |\nabla \phi|^2 + |F|^2 + U(\phi)$$
  

$$n = \text{Chern class of } W$$

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- $M_n = S_n(\Sigma)$ , Kähler
- Monopoles:  $\Sigma^3 = \mathbb{R}^3$ ,  $W = \mathbb{C}^2$  bundle
  - n = subtle
  - M<sub>n</sub> = Rat<sup>\*</sup><sub>n</sub>, hyperkähler (Kähler w.r.t. three different complex structures *I*, *J*, *K*, satisfying quaternion algebra)

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• Instantons:  $\Sigma^4 = S^4$ ,  $\mathbb{R}^4$ ,  $W = C^2$  bundle, no  $\phi$ ,

• 
$$n = c_2(W) = \int_{\Sigma} \operatorname{tr}(F \wedge F)$$

•  $M_n = \{\text{self-dual connexions}\}$  (meaning \*F = F), also hyperkähler

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• Calorons: 
$$\Sigma^4 = \mathbf{S}^1 \times \mathbb{R}^3, T^2 \times \mathbb{R}^2, \dots$$

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- Monopoles = translation invariant instantons
- Vortices = SO(3) invariant instantons

- Volume, diameter of M<sub>n</sub>?
- Curvature properties?
- Periodic geodesics?
- Ergodicity?
- Symplectic geometry of (M<sub>n</sub>, ω)?

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## Open questions: Quantization



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### **Open questions: Quantization**



• Wavefunction  $\Psi : \mathbb{R} \times M_n \to \mathbb{C}$ 

$$i\frac{\partial\Psi}{\partial t} = \frac{1}{2}\Delta\Psi$$

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Spectral geometry of  $M_n$ 

# Open questions: Validity

Geodesic approximation based on physical intuition. Can we prove it works, i.e. rigorously bound errors?

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### **Open questions: Validity**

- Geodesic approximation based on physical intuition. Can we prove it works, i.e. rigorously bound errors?
- Proto-theorem: consider one-parameter family of IVPs for field equation with initial data φ(0) = φ<sub>0</sub> ∈ M<sub>n</sub>, φ<sub>t</sub>(0) = εφ<sub>1</sub> ∈ T<sub>φ<sub>0</sub></sub>M<sub>n</sub>, ε > 0. Define time-rescaled field

 $\varphi_{\varepsilon}(\tau) = \varphi(\tau/\varepsilon).$ 

Then there exists T > 0 such that  $\varphi_{\varepsilon} : [0, T] \times \Sigma \to N$  converges uniformly, as  $\varepsilon \to 0$ , to  $\psi(\tau)$ , the geodesic in  $M_n$  with initial data  $(\varphi_0, \varphi_1)$ .

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• Proved for vortices, monopoles (Stuart),  $\mathbb{C}P^1$  lumps on  $T^2$  (JMS)

 Geodesic approximation provides a beautiful link between geometry and physics

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- Geodesic approximation provides a beautiful link between geometry and physics
- Contributions from diverse sources: Atiyah, Hitchin +...,
   Gibbons, Manton +..., Ward +..., Witten, Sen

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- Further reading: "Topological Solitons" Manton and Sutcliffe