Wave-map flow and the geometry of the space of holomorphic maps

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November 6, 2009

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$$\phi: M \to N \subset \mathbb{R}^k, \qquad E(\phi) = \frac{1}{2} \int_M |\mathrm{d}\phi|^2$$
 $(\Delta \phi)(x) \perp T_{\phi(x)} N$ 

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• E.g. 
$$N = S^2 \subset \mathbb{R}^3$$
:

 $\Delta \phi - (\phi \cdot \Delta \phi)\phi = 0$ 

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Static Heisenberg ferromagnet.

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• Conformally invariant in dimension 2.

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• dim  $\Sigma = 2$  most interesting

• Square spin lattice:  $\mathbf{S} : \mathbb{Z} \times \mathbb{Z} \to S^2$ 

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$$\frac{d\mathbf{S}_{ij}}{d\tau} = -\mathbf{S}_{ij} \times \frac{\partial H}{\partial \mathbf{S}_{ij}}, \qquad H := \sum_{i,j} J_{ij} \left[ 2 + \mathbf{S}_{ij} \cdot (\mathbf{S}_{i,j+1} + \mathbf{S}_{i+1,j}) \right]$$

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• Inhomogeneous exchange integral  $J_{ij} > 0$ : spins like to anti-align.

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•  $x = \alpha \delta$ ,  $y = \beta \delta$ ,  $t = 2\tau \delta$ 

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- $x = \alpha \delta$ ,  $y = \beta \delta$ ,  $t = 2\tau \delta$
- Assumption:

$$\begin{array}{c} \mathbf{A}_{\alpha,\beta} \\ \mathbf{B}_{\alpha,\beta} \\ J_{ij} \end{array} \right\} \quad \stackrel{\delta \to 0}{\longrightarrow} \quad \begin{cases} \mathsf{A}(x,y) \\ \mathsf{B}(x,y) \\ \mathsf{J}(x,y) \end{cases}$$

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• New fields:  $\mathbf{m} = \frac{1}{2}(\mathbf{A} + \mathbf{B}) = O(\delta), \quad \mathbf{n} = \frac{1}{2}(\mathbf{A} - \mathbf{B})$  $\mathbf{n} \times \mathbf{n}_{tt} = \mathbf{J}^2 \mathbf{n} \times (\mathbf{n}_{xx} + \mathbf{n}_{yy}) + O(\delta)$  • Leading order

$$\mathbf{n} \times \left(\mathbf{n}_{tt} - J^2 \mathbf{n}_{xx} - J^2 \mathbf{n}_{yy}\right) = \mathbf{0}$$

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• Wave map flow  $\mathbf{n} : \mathbb{R} \times \Sigma \to S^2$  where  $\Sigma$  has metric

 $g_{\Sigma} = \mathsf{J}(x,y)^{-2}(dx^2 + dy^2)$ 

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• Static problem: conformally invariant, hence independent of J

#### Leading order

$$\mathbf{n} \times \left(\mathbf{n}_{tt} - \mathbf{J}^2 \mathbf{n}_{xx} - \mathbf{J}^2 \mathbf{n}_{yy}\right) = \mathbf{0}$$

• Wave map flow  $\mathbf{n}: \mathbb{R} \times \Sigma \to S^2$  where  $\Sigma$  has metric

$$g_{\Sigma} = \mathsf{J}(x,y)^{-2}(dx^2 + dy^2)$$

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- Static problem: conformally invariant, hence independent of J
- Boundary condition:  $\Sigma = \mathbb{R}^2 \cup \{\infty\} \sim S^2$

•  $\phi: M \to N, (M, \omega^M)$  compact co-Kähler,  $(N, \omega^N)$  Kähler

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 $d^{\mathbb{C}}\phi: \mathit{T'M} \oplus \mathit{T''M} \to \mathit{T'N} \oplus \mathit{T''N}$ 

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 $d^{\mathbb{C}}\phi: T'M\oplus T''M\to T'N\oplus T''N$ 

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$$E(\phi) = \frac{1}{2} \int_{M} |\mathrm{d}\phi|^2 = \frac{1}{2} \int_{M} (|\mathrm{d}'\phi|^2 + |\mathrm{d}''\phi|^2) =: E'(\phi) + E''(\phi)$$

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- Corollary: Let  $\phi$  be holomorphic and  $\psi$  be homotopic to  $\phi.$  Then

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i.e. Holomorphic maps minimize E in their homotopy class.

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i.e. Holomorphic maps minimize *E* in their homotopy class.

•  $\phi: \Sigma \to S^2$ :  $E \ge 4\pi n$ , equality iff holomorphic



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$$u(z) = \frac{a_0 + a_1 z + \dots + a_n z^n}{b_0 + b_1 z + \dots + b_n z^n}$$

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• Boundary condition:  $b_n = 0$ . Energy localizes around zeros of u



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#### Lumps



• Boundary condition:  $b_n = 0$ . Energy localizes around zeros of u



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• Moduli space  $M_n \subset \mathbb{C}^{2n}$ 



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- Constrain  $\varphi(t)$  to  $M_n$

$$S = \int dt \{ \frac{1}{2} \int_{\Sigma} |\varphi_t|^2 - E(\varphi) \}$$

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Geodesic motion on  $(M_n, \gamma)$  where  $\gamma = L^2$  metric.

J nonconstant, bounded: M<sub>1</sub> = C × C<sup>×</sup>, foliation by smoothed out copies of Σ = C

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•  $\Sigma$  noncompact,  $\gamma$  singular (non-normalizable zero-modes)

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• Impose periodicity,  $J = \text{constant}, \Sigma = T^2$ 

#### Theorem (Sadun, JMS)

Let  $\Sigma$  be a compact Riemann surface with metric  $g_{\Sigma}$ ,  $M_n$  be the (smooth locus) of the space of degree *n* holomorphic maps  $\Sigma \to S^2$ , and  $\gamma$  be the  $L^2$  metric on  $M_n$ . Then  $(M_n, \gamma)$  is geodesically incomplete.

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**Proof:** Let  $u_0 \in M_n$ . Then  $ku_0$ , k = 1, 2, 3... is a Cauchy sequence in  $M_n$ . Its pointwise limit is discontinuous, so the sequence does not converge.  $\Box$ .

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**Proof:** Let  $u_0 \in M_n$ . Then  $ku_0$ , k = 1, 2, 3... is a Cauchy sequence in  $M_n$ . Its pointwise limit is discontinuous, so the sequence does not converge.  $\Box$ .

 As k→∞, energy localizes around zeros of u<sub>0</sub>, i.e. lumps shrink and collapse

•  $M_1 = \operatorname{Rat}_1 = PL(2, \mathbb{C}) = SO(3) \times \mathbb{R}^3 = TSO(3) = \ldots \subset \mathbb{C}P^3$ 

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• Cohomogeneity 1  $SO(3) \times SO(3)$  action

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 Kähler, finite diameter and volume, Ricci positive, unbded scalar curvature

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- Kähler, finite diameter and volume, Ricci positive, unbded scalar curvature
- Geodesic flow complicated. Generically lumps do not travel on great circles

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M<sub>2</sub>(T<sup>2</sup>) = [T<sup>2</sup> × Rat<sub>1</sub>]/[ℤ<sub>2</sub> × ℤ<sub>2</sub>], finite volume and diameter (JMS), numerics (Cova)

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•  $M_n(\mathbb{R} \times S^1)$  geodesic flow (Romao)

## Geometry of $M_n(\Sigma)$

- M<sub>2</sub>(T<sup>2</sup>) = [T<sup>2</sup> × Rat<sub>1</sub>]/[ℤ<sub>2</sub> × ℤ<sub>2</sub>], finite volume and diameter (JMS), numerics (Cova)
- $M_n(\mathbb{R} \times S^1)$  geodesic flow (Romao)
- M<sup>eq</sup><sub>n</sub>(S<sup>2</sup>) = ℝ × S<sup>1</sup> volume, total Gauss curvature, lifted geodesic flow (McGlade)

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- Spectral geometry of M<sub>1</sub>(S<sup>2</sup>): quantum dynamics of a lump on S<sup>2</sup> (Krusch-JMS)

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2nd strand: Prove conjecture (geodesic flow in  $M_n$  approximates wave map flow)

Consider one-parameter family of Cauchy problems for wave map flow  $\mathbb{R} \times \Sigma \to S^2$ :

 $\varphi(0) = \varphi_0, \qquad \varphi_t(0) = \varepsilon \varphi_1$ 

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where  $\phi_0 \in M_n$ ,  $\phi_1 \in T_{\phi_0}M_n$  and  $\varepsilon > 0$ .

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There exist T > 0 and  $\varepsilon_* > 0$  (depending on  $(\phi_0, \phi_1)$ ) such that, for all  $\varepsilon \in (0, \varepsilon_*]$ , Cauchy problem has a unique solution for  $t \in [0, T/\varepsilon]$ .

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 $\varphi_{\varepsilon}: [0,T] \times \Sigma \to S^2, \qquad \varphi_{\varepsilon}(\tau,x) = \varphi(\tau/\varepsilon,x)$ 

converges uniformly to  $\psi : [0, T] \times \Sigma \to S^2$ , the geodesic in  $M_n$  with the same initial data.

We'll sketch the proof in the case  $\Sigma = T^2$ . Ingredients:

● Wave map eqn for  $\phi \leftrightarrow$  coupled ODE/PDE system for  $\phi = \psi + \epsilon^2 Y$ 

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Ocercivity of the Hessian (and "higher" Hessian)

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- Ocercivity of the Hessian (and "higher" Hessian)
- Energy estimates for Y(t)

We'll sketch the proof in the case  $\Sigma = T^2$ . Ingredients:

- Wave map eqn for  $φ \leftrightarrow$  coupled ODE/PDE system for  $φ = ψ + ε^2 Y$
- Short time existence and uniqueness theorem for this system (in a suitable Sobolev space)

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- Ocercivity of the Hessian (and "higher" Hessian)
- Energy estimates for Y(t)
- A priori bound
• dim<sub> $\mathbb{R}$ </sub>M<sub>n</sub>(T<sup>2</sup>) = 4n (n  $\geq$  2)



- dim<sub> $\mathbb{R}$ </sub>M<sub>n</sub>(T<sup>2</sup>) = 4n (n  $\geq$  2)
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- Choose and fix initial data  $\phi_0 \in M_n$ ,  $\phi_1 \in T_{\phi_0}M_n$ .
- Choose and fix real local coords q : ℝ<sup>4n</sup> ⊃ U → M<sub>n</sub> Denote by ψ(q) the h-map corresponding to q.
   Convenient to demand that φ<sub>0</sub> = ψ(0) and U = ℝ<sup>4n</sup>.

• Wave map equation

$$\varphi_{tt} - \varphi_{xx} - \varphi_{yy} + (|\varphi_t|^2 - |\varphi_x|^2 - |\varphi_y|^2)\varphi = 0$$

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- Slow time  $\tau = \varepsilon t$  (book-keeping device)
- Decompose  $\varphi(t) = \psi(q(\tau)) + \varepsilon^2 Y(t)$ .
  - Section: map  $Z: \Sigma \to \mathbb{R}^3$

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Y is **not** a tangent section (but it's close):

$$|\psi|^2 = |\phi|^2 = 1 \quad \Rightarrow \quad \psi \cdot Y = -\frac{1}{2}\varepsilon^2 |Y|^2$$

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$$\langle Y, Z \rangle_{L^2} = 0, \qquad \forall Z \in T_{\psi(q)} \mathsf{M}_n.$$

 $Y_{tt} + L_{\Psi}Y = k + \varepsilon j$ 

where

$$L_{\Psi}Y = -\Delta Y - (|\psi_x|^2 + |\psi_y|^2)Y - 2(\psi_x \cdot Y_x + \psi_y \cdot Y_y)\psi + 2\{(\psi \cdot Y)\Delta \psi + (\psi \cdot Y)_x\psi_x + (\psi \cdot Y)_y\psi_y\}$$

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+2{ $(\psi \cdot Y)\Delta\psi + (\psi \cdot Y)_{x}\psi_{x} + (\psi \cdot Y)_{y}\psi_{y}$ }  
$$k = -\psi_{\tau\tau} - |\psi_{\tau}|^{2}\psi$$
  
$$j = j(\psi, \psi_{\tau}, Y, Y_{t}, Y_{x}, Y_{y})$$

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$$\left. \frac{d^2 E(\psi_s)}{ds^2} \right|_{s=0} = \langle Y, J_{\psi} Y \rangle, \qquad (Y = \partial_s \psi_s|_{s=0})$$

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- $T_{\Psi}M_n = \ker J_{\Psi}$
- ker  $L_{\psi} = \ker J_{\psi} \oplus \langle \psi \rangle$

# Evolution of $q(\tau)$

• Recall L<sup>2</sup> orthogonality constraint

$$\langle Y, \frac{\partial \Psi}{\partial q^i} 
angle = 0, \qquad i = 1, 2, \dots, 4n$$
  
since  $\frac{\partial \Psi}{\partial q^i}$  span ker  $J_{\Psi}$ 

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• Differentiate w.r.t. t twice

$$\langle Y_{tt}, \frac{\partial \Psi}{\partial q^i} \rangle = O(\varepsilon)$$

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 $\langle -LY + k, \frac{\partial \Psi}{\partial q^i} \rangle = O(\varepsilon)$ 

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Geodesic flow (with  $O(\varepsilon)$  correction).

# Summary: the ODE/PDE system

$$Y_{tt} + LY = k + \varepsilon j$$
  
$$q_{\tau\tau}^{i} + \Gamma(q)_{jk}^{i} q_{\tau}^{j} q_{\tau}^{k} = \varepsilon f^{i}(q, q_{\tau}, Y, Y_{t}, \varepsilon)$$

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#### Short time existence theorem

There exist  $\varepsilon_*$ , T > 0, depending only on  $\Gamma$ , such that, for all  $\varepsilon \in (0, \varepsilon_*]$  and any initial data

# $\|Y(0)\|_3^2 + \|Y_t(0)\|_2^2 + |q(0)|^2 + |q_t(0)|^2 \le \Gamma^2$

the system has a unique solution

 $(Y,q) \in C^0([0,T], H^3 \oplus \mathbb{R}^{4n}) \cap \cdots \cap C^3([0,T], H^0 \oplus \mathbb{R}^{4n})$ 

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Proof: Picard's method.

• Coercivity: show that  $\langle LY, LLY \rangle$  controls  $||Y||_3^2$ , and  $||LY_t||_0^2$  controls  $||Y_t||_2^2$  for  $Y \perp \ker J$ 

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grows slowly (would be conserved if  $\psi$  were constant,  $\varepsilon = 0$ ).

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- "Precise conjecture" follows:
- $q(\tau) 
  ightarrow q_0(\tau)$  uniformly on [0, T]
- $||Y||_3$  remains bounded for  $t \in [0, T/\epsilon]$ , and

 $\|Y\|_{C^0} \le c \|Y\|_2 \le c \|Y\|_3,$ 

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so  $\phi_{\epsilon}(\tau)$  converges uniformly on [0, T] to  $\psi(q_0(\tau))$ 

Proved for Σ = T<sup>2</sup>, but argument immediately generalizes to any compact Riemann surface

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- Can't do much better: M<sub>n</sub> incomplete!
- *T* depends on how close  $\varphi_0$  is to  $\partial M_n$ .
## Concluding remarks

- Proved for  $\Sigma = T^2$ , but argument immediately generalizes to any compact Riemann surface
- Loosely: the geodesic approximation "works" for times of order  $1/\epsilon$  when the initial velocities are of order  $\epsilon$
- Can't do much better: M<sub>n</sub> incomplete!
- *T* depends on how close  $\varphi_0$  is to  $\partial M_n$ .
- Geodesic approx. certainly fails very close to blow-up (Numerics: Linhart-Sadun, Bizoń-Chmaj-Tabor, Analysis: Rodnianski and Sterbenz)

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