

# The $L^2$ geometry of moduli spaces of $\mathbb{P}^1$ vortices

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- $X = S^2$ ,  $G = U(1)$ ,  $\mu(\mathbf{n}) = \mathbf{e} \cdot \mathbf{n} - \tau$
- $P$ , principal  $G$ -bundle over  $\Sigma$
- Connexion  $A$  on  $P$ , section  $\mathbf{n}$  of  $P^X$

$$E(\mathbf{n}, A) = \frac{1}{2} \int_{\Sigma} (|d_A \mathbf{n}|^2 + |F_A|^2 + |\mu(\mathbf{n})|^2)$$

- Vortex equations

$$(V1): \quad \bar{\partial}_A \mathbf{n} = 0, \quad (V2): \quad *F_A = \sharp \mu(\mathbf{n})$$

- General fact: solutions (if they exist) minimize  $E$  in their homotopy class

- Canonical sections  $\mathbf{n}_+(p) = \mathbf{e}$ ,  $\mathbf{n}_-(p) = -\mathbf{e}$
- Topological invariants:  $k_{\pm} = \#(\mathbf{n}(\Sigma), \mathbf{n}_{\pm}(\Sigma))$   
Assume  $k_+ \geq k_- \geq 0$
- $\deg P = k_+ - k_-$
- $E(\mathbf{n}, A) \geq 2\pi(1 - \tau)k_+ + 2\pi(1 + \tau)k_-$   
with equality iff (V1), (V2)
  - (+)-vortices located at  $\mathbf{n}^{-1}(\mathbf{e})$ , mass  $2\pi(1 - \tau)$
  - (-)-vortices located at  $\mathbf{n}^{-1}(-\mathbf{e})$ , mass  $2\pi(1 + \tau)$
- “Bradlow” bound

$$F_A = \int_{\Sigma} (\mathbf{e} \cdot \mathbf{n} - \tau)$$
$$2\pi(k_+ - k_-) = \int_{\Sigma} (\mathbf{e} \cdot \mathbf{n} - \tau) \leq (1 - \tau)|\Sigma|$$

**Theorem** (Sibner, Sibner, Yang 2000; RG 2019):

Assume  $|\Sigma| > 2\pi(k_+ - k_-)/(1 - \tau)$ . Then for each pair  $D_+, D_-$  of disjoint effective divisors on  $\Sigma$  there exists a unique (up to gauge) solution of (V1), (V2) with  $\mathbf{n}^{-1}(\pm \mathbf{e}) = D_{\pm}$ . This solution is smooth. If all vortex positions are distinct (all elements of  $D_{\pm}$  have multiplicity 1), the solution depends smoothly on the vortex positions also.

- Moduli space of vortices

$$M_{k_+, k_-}(\Sigma) \equiv [\text{Sym}^{k_+}(\Sigma) \times \text{Sym}^{k_-}(\Sigma)] \setminus \Delta_{fat}$$

- **Noncompact** complex mfd,  $\dim k_+ + k_-$

# Existence of vortices

- Key idea in proof: trade (V1), (V2) for “Taubes” equation

$$u = \log \left( \frac{1 - \mathbf{e} \cdot \mathbf{n}}{1 + \mathbf{e} \cdot \mathbf{n}} \right) : \Sigma \rightarrow [-\infty, \infty]$$

$$\Delta_{\Sigma} u + 2 \left( \frac{e^u - 1}{e^u + 1} - \tau \right) + 4\pi \left( \sum_{p \in D_+} \delta_p - \sum_{q \in D_-} \delta_q \right) = 0$$

- Regularize
- Green’s function:  $(\Sigma \times \Sigma) \setminus \Delta \rightarrow \mathbb{R}, (p, x) \mapsto G_p(x)$ 
  - smooth
  - $G_p(x) = G_x(p)$
  - $\int_{\Sigma} G_p = 0$
  - $\Delta_{\Sigma} G_p = \delta_p - |\Sigma|^{-1}$
  - In a nbhd of  $p$ ,  $G_p(x) = -(2\pi)^{-1} \log d(p, x) + \text{smooth}$
- $v := -4\pi \left( \sum_{p \in D_+} G_p - \sum_{q \in D_-} G_q \right)$
- $u = v + h$

$$\Delta_{\Sigma} h + F(v + h) - C_0 = 0$$

$$C_0 = 4\pi|\Sigma|^{-1}(k_+ - k_-) \geq 0$$

$$F(t) = 2(\tanh \frac{t}{2} + \tau).$$

- “Bradlow” bound implies  $F(-\infty) < C_0 < F(\infty)$
- $H^1(\Sigma) = \mathcal{X}(\Sigma) \oplus \mathbb{R}$

- Given  $\tilde{h} \in \mathcal{X}$ , there exists unique  $c \in \mathbb{R}$  s.t.

$$\int_{\Sigma} (F(v + \tilde{h} + c) - C_0) = 0$$

- $\mathcal{X} \rightarrow \mathbb{R}$ ,  $\tilde{h} \mapsto c(\tilde{h})$  is weakly cts
- Apply Leray-Schauder to  $T : \mathcal{X} \rightarrow \mathcal{X}$ ,  $\tilde{h} \mapsto H$ ,

$$\Delta_{\Sigma} H + F(v + \tilde{h} + c(\tilde{h})) - C_0 = 0$$

- Smoothness: bootstrap
- Uniqueness: max principle/monotonicity of  $F$
- Parametric smoothness: IFT

# The $L^2$ metric on $M_{k_+, k_-}(\Sigma)$

- Curve  $c(t) = (\mathbf{n}(t), A(t))$  of solns of (V1), (V2)
- Project  $(\dot{\mathbf{n}}(0), \dot{A}(0)) \perp L^2$  gauge orbit through  $(\mathbf{n}(0), A(0))$   
[Physics:  $\vec{E} = * \dot{A}$  satisfies Gauss's law]
- $\|\dot{c}(0)\|^2 := \int_{\Sigma} (|\dot{\mathbf{n}}(0)|^2 + |\dot{A}(0)|^2)$
- Defines a Riemannian metric  $g$  on  $M_{k_+, k_-}(\Sigma)$ . Very natural:
  - Kähler
  - Geodesic flow on  $(M_{k_+, k_-}(\Sigma), g) \leftrightarrow$  low energy dynamics of vortices
  - Quantum dynamics of vortices:  $i\partial_t \psi = \frac{1}{2} \Delta_g \psi$
  - Statistical mechanics of vortices:  $\text{Vol}(M_{k_+, k_-}, g)$  in large  $k_+, k_-$  limit

# The $L^2$ metric on $M_{k_+, k_-}(\Sigma)$

- How does one compute  $g$  in practice?
  - Cover (almost all)  $\Sigma$  with a coordinate patch
  - Consider collection of time dependent vortex trajectories  $(z_1(t), \dots, z_{k_- + k_+}(t))$
  - Construct  $(\mathbf{n}(t), A(t))$ , project
  - Compute  $\int_{\Sigma} (|\dot{\mathbf{n}}(0)|^2 + |\dot{A}(0)|^2)$ .
  - Absolutely hopeless
- Amazing fact:  $\int_{\Sigma}$  localizes around vortex positions



# The $L^2$ metric on $M_{k_+, k_-}(\Sigma)$

- $\pi : \Sigma^{k_+ + k_-} \setminus C \rightarrow M_{k_+, k_-}(\Sigma)$

**Thm**(NR, JMS 2018)

$$\pi^* \omega_{L^2} = 2\pi(1 - \tau) \sum_{r=1}^{k_+} p r_r^* \omega_\Sigma + 2\pi(1 + \tau) \sum_{r=k_++1}^{k_++k_-} p r_r^* \omega_\Sigma + i\pi \partial b$$

where  $b = \sum_r b_r d\bar{z}_r$  and  $b_r(z_1, \dots, z_{k_-+k_+})$  are defined by

$$\pm u(z) = \log |z - z_r|^2 + a_r + \frac{b_r}{2}(\bar{z} - \bar{z}_r) + \frac{\bar{b}_r}{2}(z - z_r) + \dots$$

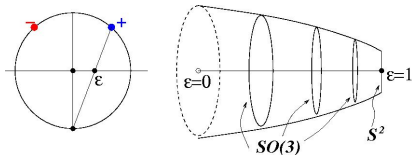
- “Strachan-Samols” localization: gives  $g$  almost explicitly on complement of coincidence set
- Moral: to compute  $g$  we only need to know how  $\mathbf{e} \cdot \mathbf{n}$  (equiv.  $u$ , equiv.  $h$ ) behaves in a nbhd of vortex positions  
In particular: how does  $dh$  at  $z_r$  depend on  $(z_1, \dots, z_{k_-+k_+})$ ?

# Volume of $M_{1,1}(S_R^2)$

- $M_{1,1}(\Sigma) = \Sigma \times \Sigma \setminus \Delta$

$$\omega_{L^2} = 2\pi(1 - \tau)\omega_{\Sigma}^+ + 2\pi(1 + \tau)\omega_{\Sigma}^- + i\pi\partial b$$

$$b = b_+ d\bar{z}_+ + b_- d\bar{z}_-$$



**Thm**(NR, JMS 2018, RG, JMS 2019)

$$\text{Vol}(M_{1,1}(S_R^2)) = 2\pi(1 - \tau)|S_R^2| \times 2\pi(1 + \tau)|S_R^2|$$

*Proof:* symmetry,

$$\omega_{L^2} = A'(\epsilon)d\epsilon \wedge \sigma_3 + A(\epsilon)\sigma_1 \wedge \sigma_2 + \frac{c}{1 + \epsilon^2} \left( \frac{1 - \epsilon^2}{1 + \epsilon^2} d\epsilon \wedge \sigma_1 - \epsilon \sigma_2 \wedge \sigma_3 \right)$$

- $\text{Vol} = 4\pi^2 \lim_{\epsilon \rightarrow 0} A(\epsilon) - c^2 \pi^2$

# Volume of $M_{1,1}(S_R^2)$

- localization formula  $\Rightarrow A(\varepsilon)$  in terms of  $b_+(\varepsilon, -\varepsilon)$
- it's enough to show  $|dh_\varepsilon| \leq C$  as  $\varepsilon \rightarrow 0$ .
- So we need to understand solution  $h_\varepsilon$  of

$$\Delta_\Sigma h_\varepsilon + F(v_\varepsilon + h_\varepsilon) = 0$$

where  $v_\varepsilon = -2\pi(G_\varepsilon - G_{-\varepsilon})$  in coalescing limit,  $\varepsilon \rightarrow 0$ .

- Naively:  $v_\varepsilon \rightarrow 0$  pointwise, suggests  $h_\varepsilon \rightarrow h_0$  where

$$\Delta_\Sigma h_0 + F(h_0) = 0$$

Max principle, monotonicity  $\Rightarrow h_0 = c_* = F^{-1}(0)$

**Thm**(RG, JMS 2019) On any  $\Sigma$ ,  $\|h_\varepsilon - c_*\|_{C^1} \rightarrow 0$  as  $\varepsilon \rightarrow 0$ .

*Proof:*  $h_\varepsilon = \tilde{h}_\varepsilon + c_\varepsilon$

- $\Delta_\Sigma \tilde{h}_\varepsilon = -F(v_\varepsilon + \tilde{h}_\varepsilon + c_\varepsilon) \Rightarrow \|\Delta_\Sigma \tilde{h}_\varepsilon\|_{L^2} \leq C$
- SEE  $\Rightarrow \|\tilde{h}_\varepsilon\|_{H^2} \leq C$
- Sobolev  $\Rightarrow \|\tilde{h}_\varepsilon\|_{C^0} \leq C$
- $\int_\Sigma F(v_\varepsilon + \tilde{h}_\varepsilon + c_\varepsilon) = 0 \Rightarrow |c_\varepsilon| \leq C$
- Alaoglu/Rellich-Kondrachov/Bolzano-Weierstrass:  
 $\tilde{h}_\varepsilon \rightharpoonup h'$  in  $H^2$ ,  $\tilde{h}_\varepsilon \rightarrow h'$  in  $H^1$ ,  $c_\varepsilon \rightarrow c'$
- $v_\varepsilon \rightarrow 0$  in  $L^2$ , MVT  $\Rightarrow F(v_\varepsilon + \tilde{h}_\varepsilon + c_\varepsilon) \rightarrow F(h' + c')$
- $h'$  weak soln of  $\Delta_\Sigma h' + F(h' + c') = 0$
- Max principle  $\Rightarrow h' = 0$ ,  $c' = c_*$

# Coalescing vortices

**Thm**(RG, JMS 2019) On any  $\Sigma$ ,  $\|h_\varepsilon - c_*\|_{C^1} \rightarrow 0$  as  $\varepsilon \rightarrow 0$ .

*Proof cont:*

- $\tilde{h}_\varepsilon \rightarrow 0$  in  $H^1$ ,  $c_\varepsilon \rightarrow c_*$ . MVT

$$|F(v_\varepsilon + \tilde{h}_\varepsilon + c_\varepsilon) - F(c_*)| \leq 2|v_\varepsilon + \tilde{h}_\varepsilon + c_\varepsilon - c_*|$$

$$\Rightarrow \|\Delta_\Sigma \tilde{h}_\varepsilon\|_{L^2} \leq 2(\|v_\varepsilon\|_{L^2} + \|\tilde{h}_\varepsilon\|_{L^2} + |c_\varepsilon - c_*|) \rightarrow 0$$

- SEE  $\|\tilde{h}_\varepsilon\|_{H^2} \rightarrow 0$ , Sobolev  $\|h_\varepsilon\|_{C^0} \rightarrow 0$
- Calderon-Zygmund:  $\|f\|_{L^2_p} \leq C(\|\Delta_\Sigma f\|_{L^p} + \|f\|_{L^p})$
- LDCT  $\Rightarrow$  for all  $p > 2$

$$\|h_\varepsilon\|_{L^2_p} \rightarrow 0$$

- Sobolev  $\Rightarrow \|\tilde{h}_\varepsilon\|_{C^1} \rightarrow 0$ .  $\square$

- $p_\varepsilon := \partial_\varepsilon h_\varepsilon$

$$\Delta_\Sigma p_\varepsilon + F'(v_\varepsilon + h_\varepsilon)p_\varepsilon = -F'(v_\varepsilon + h_\varepsilon)\partial_\varepsilon v_\varepsilon$$

Lax-Milgram gives estimate for  $\|p_\varepsilon\|_{H^1} \dots$

- **Thm**(RG, JMS 2019)  $\|p_\varepsilon\|_{H^3} \leq C/\varepsilon$
- **Cor**(NR, JMS 2018, RG, JMS 2019)  
 $(M_{1,1}(\Sigma), g)$  is geodesically incomplete

# Compactification of $M_{k_+, k_-}(\Sigma)$

$$[\text{Sym}^{k_+}(\Sigma) \times \text{Sym}^{k_-}(\Sigma)] \setminus \Delta_{\text{fat}} \hookrightarrow \text{Sym}^{k_+}(\Sigma) \times \text{Sym}^{k_-}(\Sigma)$$

- Can we extend  $g$  smoothly to R.H. mfd? No!
- Identify R.H. mfd as moduli space of vortices in a **linear** gauged sigma model
- $X = \mathbb{C}^2$ ,  $G = T^2$ ,  $g_{T^2} = d\theta_1^2/e^2 + d\theta_2^2$

$$(\lambda_1, \lambda_2) : (X_+, X_-) \mapsto (\lambda_1 \lambda_2 X_+, \lambda_1 X_-)$$

- Moment map

$$\begin{aligned}\mu_1(X_+, X_-) &= \frac{1}{2}(4 - |X_+|^2 - |X_-|^2) \\ \mu_2(X_+, X_-) &= \frac{1}{2}(2 - 2\tau - |X_+|^2)\end{aligned}$$

# Compactification of $M_{k_+, k_-}(\Sigma)$

- Vortex equations

$$\overline{\partial^A} \varphi_{\pm} = 0 \quad (1)$$

$$*F_{A_1} = \frac{e^2}{2}(4 - |\varphi_+|^2 - |\varphi_-|^2) \quad (2)$$

$$*F_{A_2} = 1 - \tau - \frac{1}{2}|\varphi_+|^2 \quad (3)$$

- Provided  $k_+ \geq k_- > \max\{0, 2\text{genus}(\Sigma) - 2\}$  and a Bradlow bound is satisfied

$$M_{k_+, k_-}^{C^2, e}(\Sigma) \equiv \text{Sym}^{k_+}(\Sigma) \times \text{Sym}^{k_-}(\Sigma)$$

- One-parameter family of metrics  $g_e$  on  $\text{Sym}^{k_+}(\Sigma) \times \text{Sym}^{k_-}(\Sigma)$
- $\iota^* g_e$ : one-parameter family of metrics on  $M_{k_+, k_-}(\Sigma)$



# Compactification of $M_{k_+, k_-}(\Sigma)$

**Conjecture**(NR, JMS 2018)  $\iota^* g_e$  converges uniformly to  $g$  (the  $L^2$  metric on  $M_{k_+, k_-}(\Sigma)$ ) as  $e \rightarrow \infty$ .

- Motivation: forgetful map

$T : (\varphi_+, \varphi_-, A_1, A_2) \mapsto ([\varphi_+ : \varphi_-], A_2)$  globalizes

$$\Gamma(P^{\mathbb{C}^2}) \times \mathcal{A}(P) \rightarrow \Gamma(P_2^{\mathbb{P}^1}) \times \mathcal{A}(P_2)$$

Formally a Riemannian submersion

- For fixed pair of (disjoint) divisors, apply to  $(\varphi_+^e, \varphi_-^e, A_1^e, A_2^e)$  solution of vortex equations
  - $\bar{\partial}_{A_2^e}[\varphi_+^e : \varphi_-^e] = 0$  automatically — solves (V1)
  - Expect  $|\varphi_+^e|^2 + |\varphi_-^e|^2 = 4 + O(e^{-2})$
  - Then  $*F_{A_2^e} = \frac{|\varphi_+^e|^2 - |\varphi_-^e|^2}{|\varphi_+^e|^2 + |\varphi_-^e|^2} - \tau + O(e^{-2})$
  - So  $([\varphi_+^e : \varphi_-^e], A_2)$  solves (V2) up to an error of order  $e^{-2}$
- Similar conjecture for **ungauged** maps  $\Sigma \rightarrow \mathbb{P}^{n-1}$  and  $U(1)$  vortices with  $X = \mathbb{C}^n$  proved by Liu

# Testing the conjecture

- Conjecture implies

$$\text{Vol}(M_{k_+, k_-}(\Sigma), g) = \lim_{e \rightarrow \infty} \text{Vol}(M_{k_+, k_-}^{\mathbb{C}^2, e}(\Sigma, g_e))$$

- Using ideas of Baptista, can write down Kähler class of  $M^{\mathbb{C}^2}$  exactly
- Computing

$$\text{Vol}(M^{\mathbb{C}^2}, g_e) = \int_{M^{\mathbb{C}^2}} \frac{\omega_e^{k_+ + k_-}}{(k_+ + k_-)!}$$

reduces to an exercise in understanding the cohomology ring of  $M^{\mathbb{C}^2} = \text{Sym}^{k_+}(\Sigma) \times \text{Sym}^{k_-}(\Sigma)$

**Thm** (NR, JMS 2019)

$$\text{Vol} \left( M_{k_+, k_-}^{\mathbb{C}^2, e}(\Sigma) \right) = \sum_{\ell=0}^g \frac{g!(g-\ell)!}{(-1)^{\ell\ell!}} \prod_{\sigma=\pm} \sum_{j_\sigma=\ell}^g \frac{(2\pi)^{2\ell} J_\sigma^{k_\sigma-j_\sigma} K_\sigma^{j_\sigma-\ell}}{(j_\sigma-\ell)!(g-j_\sigma)!(k_\sigma-j_\sigma)!}.$$

where

$$J_+ := 2\pi(1-\tau)|\Sigma| - 4\pi^2(k_+ - k_-),$$

$$J_- := 2\pi(1+\tau)|\Sigma| - 4\pi^2 e^{-2} k_- + 4\pi^2(k_+ - k_-),$$

$$K_+ := 4\pi^2,$$

$$K_- := 4\pi^2(1 + e^{-2}).$$

# Testing the conjecture

- So the conjecture implies

$$\text{Vol}(M_{k_+, k_-}(\Sigma)) = \sum_{\ell=0}^g \frac{g!(g-\ell)!}{(-1)^\ell \ell!} \prod_{\sigma=\pm} \sum_{j_\sigma=\ell}^g \frac{(2\pi)^{2\ell} J_\sigma^{k_\sigma - j_\sigma} K_\sigma^{j_\sigma - \ell}}{(j_\sigma - \ell)!(g - j_\sigma)!(k_\sigma - j_\sigma)!}$$

where

$$J_+ := 2\pi(1 - \tau)|\Sigma| - 4\pi^2(k_+ - k_-),$$

$$J_- := 2\pi(1 + \tau)|\Sigma| + 4\pi^2(k_+ - k_-),$$

$$K_\pm := 4\pi^2.$$

- This is consistent with  $\text{Vol}(M_{1,1}(S_R^2))$ .
- Also checked (RG)  $M_{k_+, 0}(S_R^2)$ ,  $M_{0, k_-}(S_R^2)$
- No cases with  $\text{genus}(g) > 0$  have been checked. Note that the conjecture implies

$$\text{Vol}(M_{1,1}(\Sigma)) \neq 2\pi(1 - \tau)|\Sigma| \times 2\pi(1 + \tau)|\Sigma|$$

when  $g = 1$

# Testing the conjecture

- Can also compute Einstein-Hilbert action of  $M_{k_+,k_-}^{\mathbb{C}^2,e}(\Sigma)$  explicitly.
- Get conjectures for EH of  $M_{k_+,k_-}(\Sigma)$ .
- **None** of these have been checked.